



$O(N^2)$ Universal Antisymmetry in Fermionic Neural Networks

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Slater Determinant

- Not universal due to single-electron dependence

$$\psi_{\text{single}}(\mathbf{x}) = \begin{vmatrix} \phi_1(\mathbf{x}_1) & \cdots & \phi_1(\mathbf{x}_N) \\ \vdots & & \vdots \\ \phi_N(\mathbf{x}_1) & \cdots & \phi_N(\mathbf{x}_N) \end{vmatrix} = \det[\phi_i(\mathbf{x}_j)].$$

Fermionic Neural Network (FermiNet)

- Single-electron dependence to **multi-electron** dependence

$$\psi_{\text{Fermi}}(x) = \det [\phi_i(x_j; \{x_{\setminus j}\})].$$

Lemma 1. (*Universality by Pfau et al. (2020)*) For any antisymmetric function $\Psi(x)$, there exist multi-electron functions ϕ_1, \dots, ϕ_N , such that $\forall x$, there is $\psi_{\text{Fermi}}(x) = \Psi(x)$.

Reducing $O(N^3)$ Computation to $O(N^2)$

- Construction of the pair-wise Ansatz

$$\psi_{\text{pair}}(x) = \prod_{1 \leq i < j \leq N} \mathcal{A} \circ F(x_i, x_j; \{x \setminus \{i, j\}\}),$$

Lemma 2. (*Antisymmetry*) $\psi_{\text{pair}}(x)$ is antisymmetric under the permutation of any two elements x_m and x_n in x .

Instantiation from Han et al. (2019)

- Constructed via a two-electron function

$$\psi_{\text{Han}}(\mathbf{x}) = \phi_C(\mathbf{x}) \cdot \prod_{1 \leq i < j \leq N} (\phi_B(\mathbf{x}_j, \mathbf{x}_i) - \phi_B(\mathbf{x}_i, \mathbf{x}_j)).$$

Theorem 1. (*Discontinuity*) *There exist ground-state wave-functions $\Psi^g(\mathbf{x})$, such that if there is $\forall \mathbf{x}, \psi_{\text{Han}}(\mathbf{x}) = \Psi^g(\mathbf{x})$, then either ϕ_B or ϕ_C must be discontinuous on \mathbb{R}^{dN} .*

Achieving Universality with Continuous Constructions

- Multi-electron function is critical!

$$\psi'_{\text{pair}}(x) = \prod_{1 \leq i < j \leq N} (\phi_B(x_j; \{x_{\setminus j}\}) - \phi_B(x_i; \{x_{\setminus i}\})).$$

Theorem 2. (*Continuous universality*) For any ground-state wavefunction $\Psi^g(x)$, there exist a continuous multi-electron function ϕ_B , such that $\forall x, \psi'_{\text{pair}}(x) = \Psi^g(x)$.

Connection with FermiNet

- Our pair-wise construction is a special case (Vandermonde matrix)

$$\psi'_{\text{pair}}(\mathbf{x}) = \begin{vmatrix} 1 & \cdots & 1 \\ \phi_B(\mathbf{x}_1; \{\mathbf{x}_{\setminus 1}\}) & \cdots & \phi_B(\mathbf{x}_N; \{\mathbf{x}_{\setminus N}\}) \\ \phi_B^2(\mathbf{x}_1; \{\mathbf{x}_{\setminus 1}\}) & \cdots & \phi_B^2(\mathbf{x}_N; \{\mathbf{x}_{\setminus N}\}) \\ \vdots & \vdots & \vdots \\ \phi_B^{N-1}(\mathbf{x}_1; \{\mathbf{x}_{\setminus 1}\}) & \cdots & \phi_B^{N-1}(\mathbf{x}_N; \{\mathbf{x}_{\setminus N}\}) \end{vmatrix},$$

Extension

- Two multi-electron functions

$$\psi'_{\text{pair}}(\mathbf{x}) = \prod_{1 \leq i < j \leq N} (\phi_B(\mathbf{x}_j; \{\mathbf{x}_{\setminus j}\}) - \phi_B(\mathbf{x}_i; \{\mathbf{x}_{\setminus i}\})).$$

$$\psi''_{\text{pair}}(\mathbf{x}) = \prod_{1 \leq i < j \leq N} \begin{vmatrix} \phi_A(\mathbf{x}_i; \{\mathbf{x}_{\setminus i}\}) & \phi_A(\mathbf{x}_j; \{\mathbf{x}_{\setminus j}\}) \\ \phi_B(\mathbf{x}_i; \{\mathbf{x}_{\setminus i}\}) & \phi_B(\mathbf{x}_j; \{\mathbf{x}_{\setminus j}\}) \end{vmatrix}.$$

Initial experiments

Table 1. Ground state energy. The values of ‘Exact’ column come from [Chakravorty et al. \(1993\)](#). The values of ‘ ψ'_{pair} ’ and ‘ ψ''_{pair} ’ are averaged on the last 1,000 iterations with sampling stride of 10.

| Atom | ψ''_{pair} | ψ'_{pair} | Exact |
|------|------------------------|-----------------------|-------------|
| Li | -7.4782 | -7.4781 | -7.47806032 |
| Be | -14.6673 | -14.6664 | -14.66736 |
| B | -24.5602 | -24.4475 | -24.65391 |
| C | -37.3531 | -37.2785 | -37.8450 |
| N | -53.1855 | -53.0626 | -54.5892 |

Thank you

